

# Application of Statistical Mechanics

## Imperfect gases

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# Recommended literature

- Introduction to Statistical Mechanics by G.S. Rushbrooke, Oxford University Press, (1964 or another edition)
- Introduction to Statistical Thermodynamics by T.L. Hill, Dover Publications, New York (1986)
- Statistical Mechanics by D.A. McQuarrie, University Science Books, Sausalito (2000)

*Lecture slides:* <http://dullens.chem.ox.ac.uk/main/teaching.html>

# Imperfect gases?

- Low densities: all perfect gas

$$p = \rho k_B T$$

with  $\rho = \frac{N}{V}$

- Deviations at higher densities: Virial expansion

$$p = k_B T (\rho + B_2 \rho^2 + B_3 \rho^3 + \dots)$$

- Higher order terms: account for intermolecular interactions

Goal: develop statistical mechanics to derive expression for  $B_2$  in terms of intermolecular potential

# Content of the course

- Canonical Ensemble (recap)
- Classical Statistical Mechanics
- Second Virial Coefficient  $B_2$
- $B_2$  for imperfect gases

# Macroscopic properties of $10^{23}$ molecules?

## Thermodynamics

- Description of macroscopic systems ( $10^{23}$  molecules) in equilibrium
- Variables like volume, pressure, temperature, ....

$$p = \rho k_B T$$

$$p = k_B T (\rho + B_2 \rho^2 + B_3 \rho^3 + \dots)$$

- No interpretation or explanation at the molecular level

# Macroscopic properties of $10^{23}$ molecules?

## Quantum/Classical mechanics

- Description on a microscopic, i.e. molecular, level

$$\hat{H}\psi_i = \varepsilon_i\psi_i$$

- How to calculate macroscopic properties from molecular properties?

each (pointlike) molecule is described by 3 position and 3 momentum coordinates:  $6N$  variables required ... hmmm...

We have to use a statistical approach to tackle this

# Statistical Mechanics: independent systems

*From 2<sup>nd</sup> year stat mech:*

Boltzmann distribution

$$P_i = \frac{e^{-\beta\varepsilon_i}}{\sum_i e^{-\beta\varepsilon_i}}$$

Molecular partition function

$$q = \sum_i e^{-\beta\varepsilon_i}$$

Partition function for a particle of mass  $m$  in 3D:

$$q = \left( \frac{2\pi m k_B T}{h^2} \right)^{\frac{3}{2}} V$$

# Statistical Mechanics: independent systems

*From 2<sup>nd</sup> year stat mech:*

Example: average internal energy per particle

$$E = Nk_B T^2 \left( \frac{\partial \ln q}{\partial T} \right)_V$$

**Equipartition!**



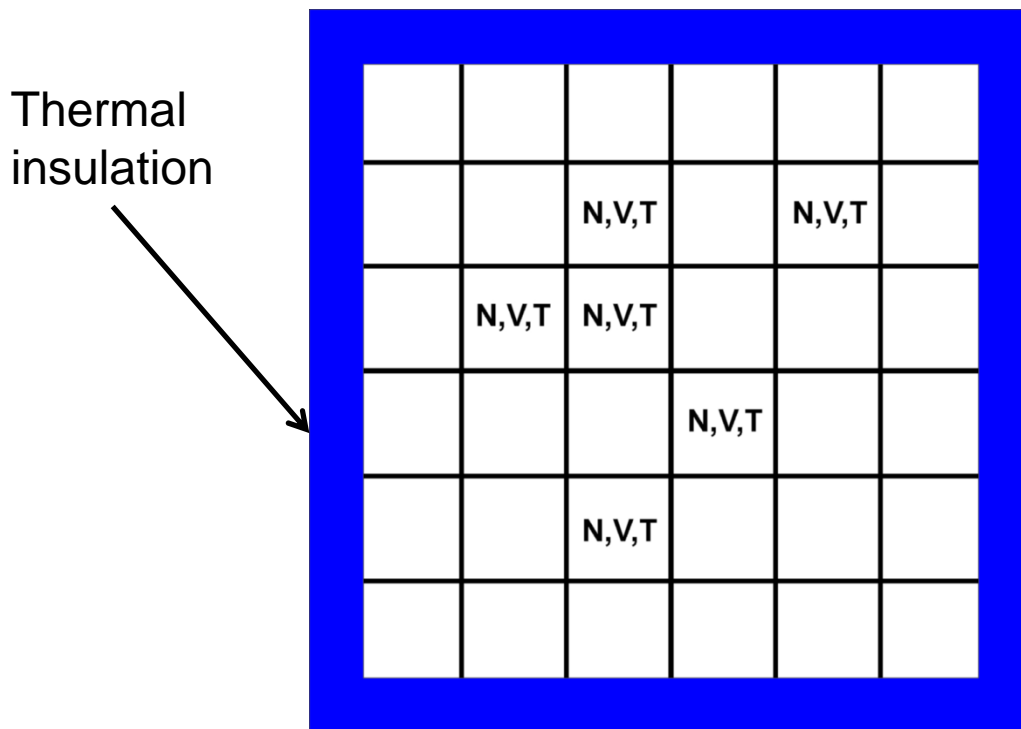
# Statistical Mechanics: interacting systems

## **Canonical Ensemble**

# Canonical Ensemble

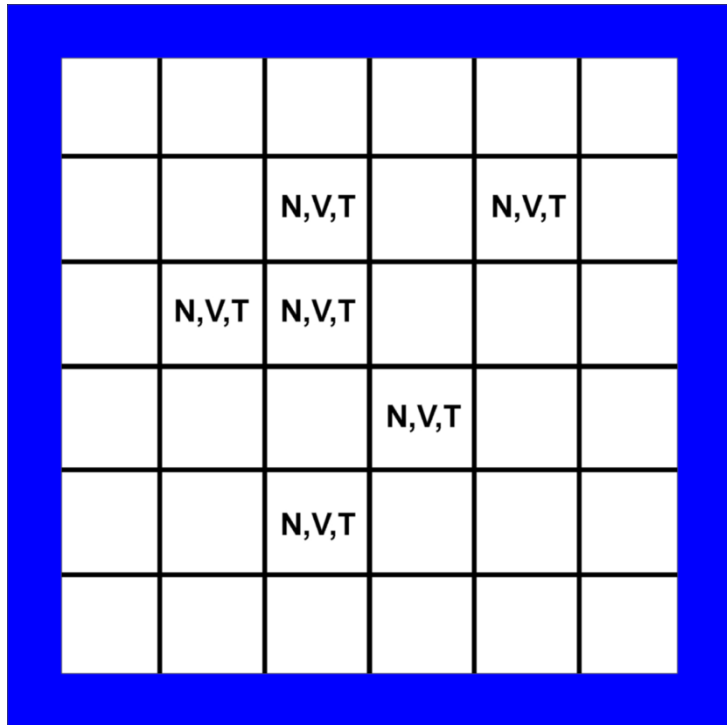
From 2<sup>nd</sup> year stat mech:

*Collection of a vary large number of systems, each replicating at a thermodynamic level the closed isothermal system of interest  
i.e. each system in ensemble has same fixed  $N, V, T$ .*



- Fixed total energy
- Fixed volume
- Fixed number of systems
- Assume:  $E_i$ 's are known

# Canonical Ensemble



		N,V,T		N,V,T	
	N,V,T	N,V,T			
			N,V,T		
		N,V,T			

Canonical distribution

$$P_i = \frac{e^{-\beta E_i}}{\sum_i e^{-\beta E_i}}$$

Canonical partition function

$$Q = \sum_i e^{-\beta E_i}$$

***Q contains all thermodynamic information about system of (in general) interacting molecules***

# Connection to thermodynamics

*From 2<sup>nd</sup> year stat mech:*

$$E = k_B T^2 \left( \frac{\partial \ln Q}{\partial T} \right)_V \qquad S = \frac{E}{T} + k_B \ln Q$$

Important for us: the equation of state ( $p$ ) in terms of  $Q$

$$p = k_B T \left( \frac{\partial \ln Q}{\partial V} \right)_{T,N}$$

# Canonical Ensemble: perfect gas

$$q = \left( \frac{2\pi m k_B T}{h^2} \right)^{\frac{3}{2}} V$$

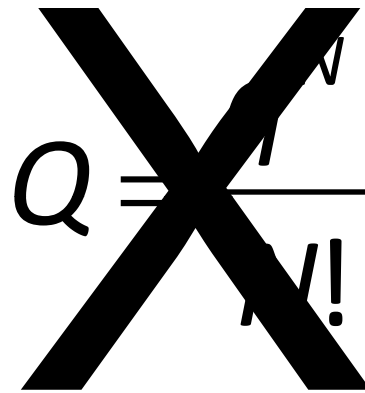
The diagram illustrates the substitution of the single-particle partition function  $q$  into the  $N$ -particle partition function  $Q$ . A light blue arrow points from the left box to the right box, and a downward-pointing arrow indicates the substitution of  $q$  into the right-hand side of the equation.

$$Q = \frac{q^N}{N!} \quad \rightarrow \quad \frac{1}{N!} \left( \frac{2\pi m k_B T}{h^2} \right)^{\frac{3N}{2}} V^N$$

Equation of state (of course):

$$p = k_B T \left( \frac{\partial \ln Q}{\partial V} \right)_{T,N} = \rho k_B T$$

# Canonical Ensemble: imperfect gas

$$Q = \frac{1}{N!} \int \dots \int e^{-\beta H} \mathcal{D}\mathbf{r}^N$$


**The trouble is finding the canonical partition function for interacting systems\***

\* This is actually the point of this lecture course

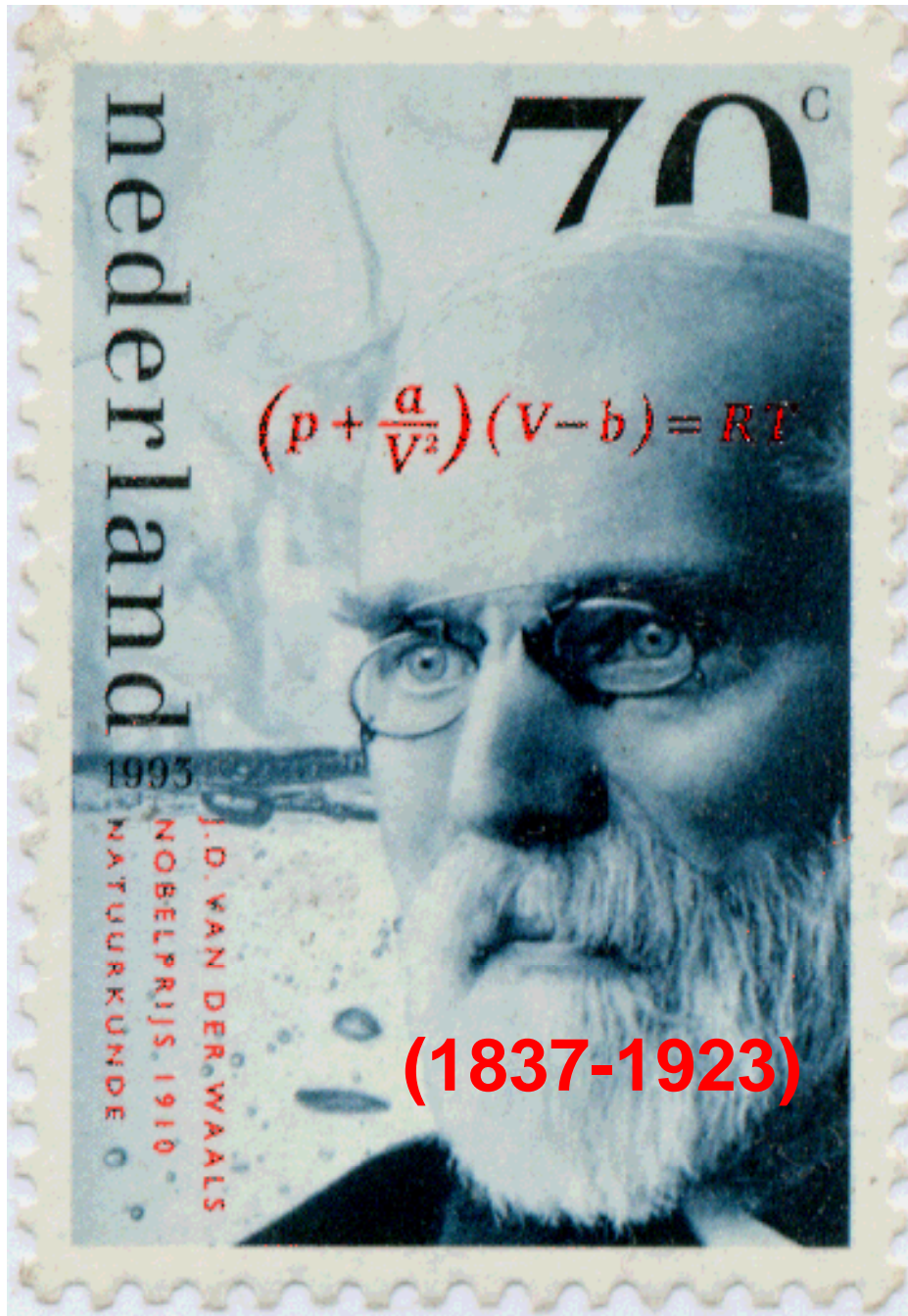
# Example of canonical partition function for an imperfect gas

$$Q = \frac{1}{N!} \left( \frac{2\pi m k_B T}{h^2} \right)^{\frac{3N}{2}} (V - Nb)^N \exp \left[ \frac{aN^2}{Vk_B T} \right]$$

Van der Waals gas

$$P = \frac{Nk_B T}{V - Nb} - a \left( \frac{N}{V} \right)^2$$

Later on: expressions for  $a$  and  $b$  (see also problem set)



# Johannes Diderik Van der Waals

Over de continuïteit van den  
gas – en vloeistof toestand.

*PhD-Thesis, Leiden, 1873*  
*The Netherlands*

Condensation requires  
*attraction* and *repulsion*



# Content of the course

Statistical mechanics to derive expressions for the canonical partition function and  $B_2$  in terms of intermolecular potential

- Canonical Ensemble (recap)
- **Classical Statistical Mechanics**
- Second Virial Coefficient  $B_2$
- $B_2$  for imperfect gases